

VII. Heat transfer coefficients

A. Review General definition of h

$$Q = h \cdot A \cdot \Delta T$$

*total rate of heat transfer* = *heat transfer coefficient*  $\cdot$  *characteristic area*  $\cdot$  *characteristic temperature difference*

B. Flow in tubes

1. Heat transfer between wall and fluid

BSLK section 14.1, 14.2

$$Q = h_{ln} (\pi DL) (T_o - T_b)_{ln} \quad \text{Eq. I}$$

$$= h_{ln} (\pi DL) \frac{(T_{o1} - T_{b1}) - (T_{o2} - T_{b2})}{\ln((T_{o1} - T_{b1}) / (T_{o2} - T_{b2}))} \quad \text{BSLK eq. 14.144}$$

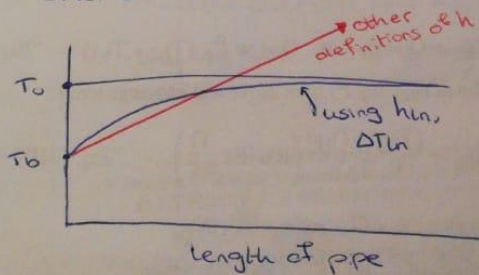
$T_{b1}$  = Fluid T entering;  $T_{b2}$  = Fluid T leaving

$T_{o1}$  = wall T entrance;  $T_{o2}$  = wall T exit

units:  $Q$  = watts

$h_{ln}$  =  $W / (m^2 K)$

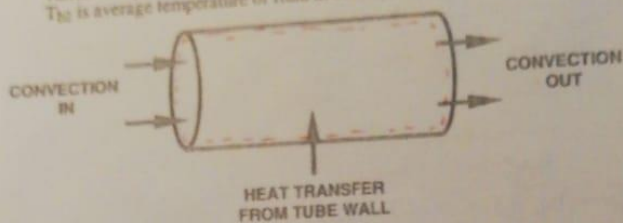
Don't pay attention to other definitions of h in BSLK sect. 14.1  
other definitions lead to unphysical situations



2. Energy balance for flow in tube

a. macroscopic balance

(Can use macroscopic energy balance because don't care about details of how  $T$  varies with position within the tube.  $T_{b1}$  is average temperature of fluid at inlet;  $T_{b2}$  is average temperature of fluid at outlet.)



Define system as fluid within pipe. Assume fluid is incompressible. Therefore average velocity  $v$  is same at inlet and outlet.

Terms in energy balance:

- convection of thermal energy in:  $(\pi R^2) v \rho \hat{C}_p T_{b1}$   
Define mass rate of flow  $w = (\pi R^2) \rho v$ ;  
then convection in =  $w \hat{C}_p T_{b1}$
- convection of thermal energy out:  $w \hat{C}_p T_{b2}$
- heat transfer from walls:  $Q$   $Q > 0 \rightarrow$  heat flows in
- No accumulation at steady state -  $T$  does not change with time at any fixed location in the pipe (though  $T$  does vary with position along the pipe).
- No heat generation within fluid.

Energy balance:

$$w \hat{C}_p T_{b1} - w \hat{C}_p T_{b2} + Q = 0; \text{ or } Q = w \hat{C}_p (T_{b2} - T_{b1}) - \text{"Eq. II"}$$

Combine with definition of  $h_{in}$  ("Eq. I") from above and rearrange terms:

$$\frac{h_{in} D}{k} \equiv Nu_{in} = \frac{(T_{b2} - T_{b1})}{(T_0 - T_{b1})_{in}} \left( Re Pr \frac{D}{4L} \right) - \text{"Eq. III"}$$

Note: IF  $T_0$  is uniform along tube ( $T_{01} = T_{02} \equiv T_0$ ), then

$$\frac{(T_{b2} - T_{b1})}{(T_0 - T_{b1})_{in}} = \ln \left( \frac{T_0 - T_{b1}}{T_0 - T_{b2}} \right) = \frac{\Delta T \text{ at inlet}}{\Delta T \text{ at outlet}}$$

- In  $\Phi E 383$ , for simplicity, we usually assume  $T_{01} = T_{02} \equiv T_0$ .
- Eq. "III" is derived from energy balance, and therefore represents conservation equation in solution of heat-transfer problems in tubes.

### 3. Correlations for $h_{in}$

As with friction factors, correlations for  $h_{in}$  relate dimensionless groups:

$$(Nu)_{in} = \frac{h_{in} D}{k} \leftarrow \begin{matrix} \text{dia. of pipe} \\ \text{thermal conductivity} \\ \text{of liquid} \end{matrix} \quad \text{"Nusselt no."}$$

$$Re = \frac{D V \rho}{\mu} = \frac{D G}{\mu}$$

$$Pr = \left( \frac{c_p \mu}{k} \right) \text{ "Prandtl no."} = \frac{\mu (c_p / \rho)}{\alpha}$$

Correlations for  $Nu_{in}$  are transport "laws" - one combines these correlations with the conservation eq., "Eq. III," to complete the solution.

The primary correlation is BSLK 14.3-2

Note:

- similarities to  $f(Re)$  chart  
 $(Nu)_{in}$  correlation with  $Re$   
 general slope
- differences with  $f(Re)$  chart
  - multiple curves at low  $Re$ , depend on  $HO$
  - only one curve in highly turbulent regime  
 (strictly, this is a chart for "smooth tubes")
- ... because correlation is only for smooth pipes
- solution in laminar region is from analytical solution to pde derived in BSLK Sect. 8.8; in turbulent region, based on experimental results  
 BSLK Sect. 10.7
- third label for the vertical axis of Fig. 14.3-2 combines transport law and conservation eq: thus can solve problem directly
- although  $[(h_{in} D/k) Re^{-1}]$  declines with increasing  $Re$  across much of chart,  $h_{in}$  increases monotonically with increasing  $Re$ . Increasing convection and turbulence increases  $h_{in}$ .

4. using the correlations for  $Nu_{ln}$

Re < 2100

- a) ~~Don't want to solve differential equation for T<sub>w</sub> with h<sub>w</sub>~~
  - use correlation below (Eq. 14.3-17, Fig. 14.3-2)
  - Fig. 14.2: "const. wall T (tube)" curve is OK for h<sub>w</sub>
  - ~~DO NOT COVERED IN PROBLEMS~~
  - ~~DO NOT COVERED IN PROBLEMS~~

- b)  $Re Pr D/L \gg 10^7$  (Or  $[(T_2 - T_b)/(T_1 - T_b)] > 0.2$ )  
(in other words, is T relatively far from equilibrium at outlet?)

- yes: then can solve in any of 3 ways:
  - use Fig 14.3-2 from BSL ~~Eq. 14.3-17~~ BSL 2 or BSLK 14.3-17
  - Eq. 14.3-17 and Eq. "III," or Eq. 3.150
  - Fig. 14.2-1 "const. wall T (tube)" curve and Eq. "III"
- no: then can use only Fig. 14.2-1 "const. wall T (tube)" & Eq. "III" ~~Eq. 3.150~~

Re > 2100

Use Fig. 14.3-2  
(or 10,000) ~~Eq. 14.3-17~~ BSL 2 or BSLK 14.3-16  
If  $Re > 20,000$  can also use Eq. ~~Eq. 14.3-17~~ and Eq. "III"

Non-cylindrical Tubes: Use "hydraulic radius" approximation  
Same approach as in friction-factor problems  
Substitute  $4R_h \approx 4(S/Z)$  for D  
Calculate  $\langle v \rangle$  for given tube shape  
Calculate h or T as for circular tube  
- VALID ONLY FOR TURBULENT FLOW

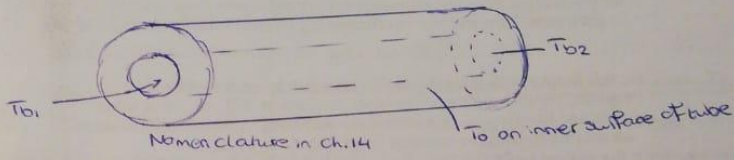
Note: for any Re, if  $T_{01} = T_{02} = T_0$ , then  
 $(T_{b2} - T_{b1}) / (T_0 - T_b)_{ln} = \ln [(T_0 - T_{b1}) / (T_0 - T_{b2})]$

in this course, we'll always assume  
 $T_{01} = T_{02} = T_0$  [unif. wall  $T_0$ ]

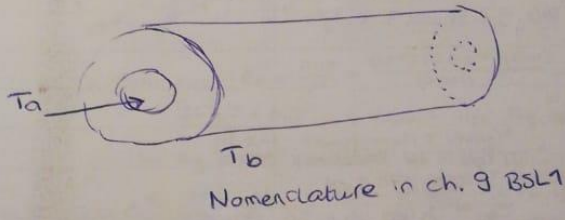
5. Additional notes

- a. Don't sweat the other correlations for  $h_{ln}$  in BSL Section 14.2, 14.3

5. Heat transfer between two fluids across wall(s)



Call  $T$  of fluid outside  $T_f$



- b.  $h_{in}$  corresponds to  $h_o$  of BSL Sect. 10.6 (BSL I Sect 7.6; BSLK Sect 10.3)
- i.e., heat-transfer coefficient for Newton's law on inner surface of tube
  - Can combine  $h_{in} = h_o$  derived with methods of BSL ch. 14 with formula for conduction through composite cylindrical layers from Section 10.6 to get overall heat-transfer coefficient  $U_o$  for heat transfer through tubes
  - (in this course, still have no way to calculate  $h_3$  on the outer tube surface; for this problem, see other sections of BSL ch. 14 or other texts and handbooks. Principles are still same: greater convection and increasing turbulence increase  $h_3$  as they do  $h_{in} = h_o$ )
- c. for  $h$  for packed beds, see BSL Sect. 14.5 - not covered in BSL 10.3-31 (BSL 10.3-31)
- d. Eq. III describes heat transfer between the inner surface of a tube (assumed to be at fixed  $T_o$ ) and fluid within the tube. What if instead have uniform  $T_f$  of fluid outside the tube? Use eq. 3.6-31 BSL 7 to get  $U_o$  with  $h_{in} = h_o$

BSLK 10.3-30  
 BSL Eqs. 10.6-29 through 31 apply, with  
 \*  $T_b$  replacing  $T_a$  in eq. 3.6-30 (fluid temp. inside tube)  
 \*  $T_p$  replacing  $T_b$  (fluid temp. outside tube)  
 calculate  $h_{in} (= h_o)$  from methods in ch. 14  
 $h_3$  must be given in problem statement  
 →  $U_o$

e. Eqs. 10.6-29 to 31 give  $U_o$ ; then how to calculate fluid T along tube?  
 Conservation equation is modified form of Eq. III:  
 From eq. 3.6-31 (BSL 7) →  $\left[ \frac{U_o D_o}{k} \right] = \ln \left( \frac{T_p - T_{b1}}{T_p - T_{b2}} \right) Re Pr \frac{D_o}{4L}$   
 $k$  thermal conductivity inside tube  
 internal diam. tube  
 $T_p$  fluid outside tube  
 $T_{b1}$  fluid entrance  
 $T_{b2}$  fluid exit

6. Example problems  
 a. BSL ex. 14.3-1

Final notes on heat transfer

The mass and energy balances in ~~PGEE333~~ and in thermodynamics are fundamentally the same, though they may look different.

In ~~PGEE333~~, we focus attention on the "Q" term, and for simplicity, leave out the following complications

- useful work output
- Phase changes. assume ~~energy~~ energy/volume is  $p\hat{c}_p T$  with constant  $p\hat{c}_p$ ; excludes heat of vaporization, accumulation

- thermodynamics addresses these two complications, without devoting so much attention to the "Q" term.

In the real world, you'll need handle situations with all factors simultaneously